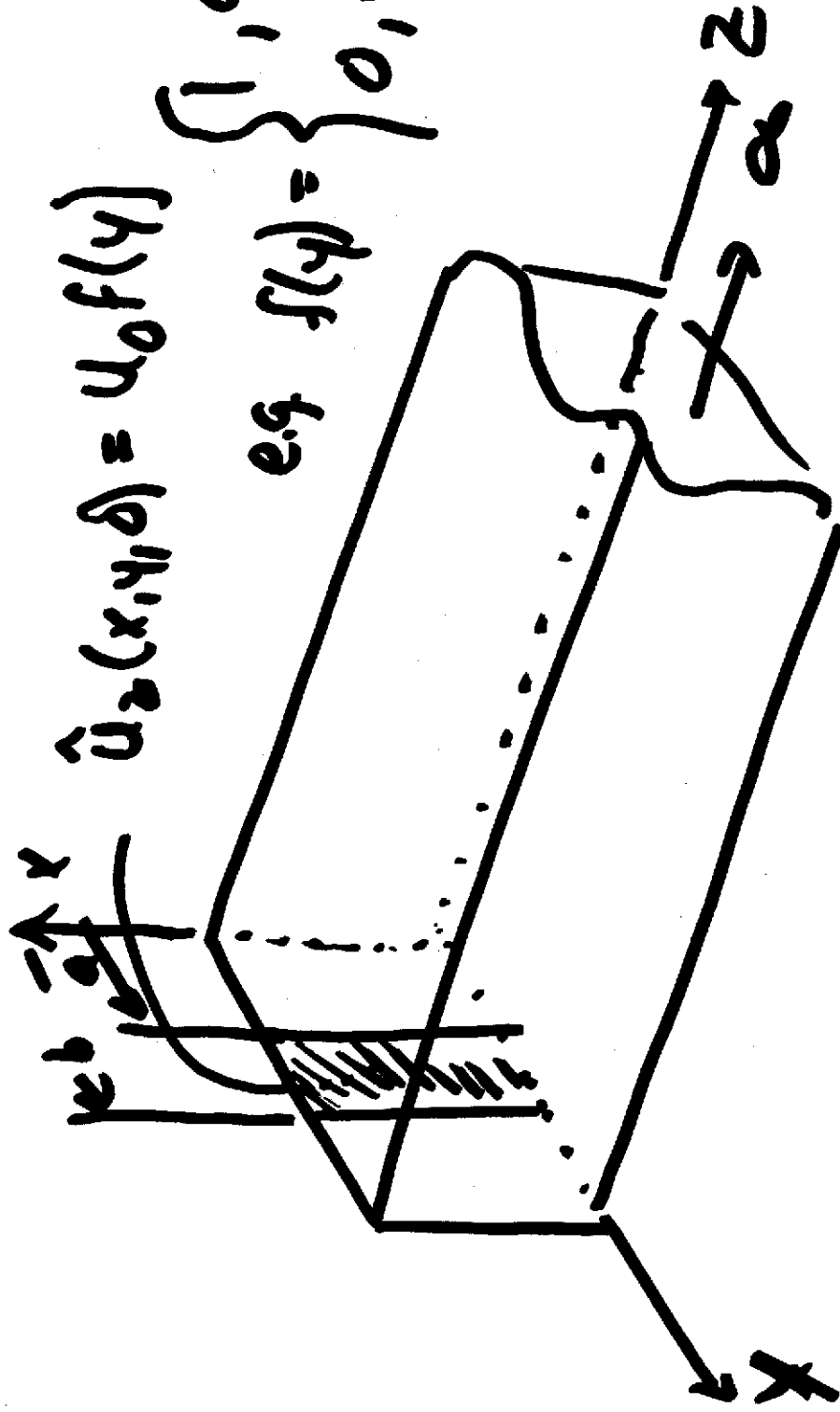


University of Idaho Consider the following problem

(2D Rectangular Waveguide)

$$\hat{u}_z(x, y, z) = u_0 f(y)$$

$$\text{e.g. } f(y) = \begin{cases} 1, & 0 \leq y \leq b \\ 0, & \text{elsewhere} \end{cases}$$



UNIVERSITY OF IDAHO Then, it is reasonable to expect that the acoustic pressure will not depend on the x

$$\text{Coordinates} \left(\begin{array}{l} e^{j\omega t} \\ \sum_{l=0}^{\infty} \sum_{m=0}^{\infty} \hat{A}_{lm} \cos\left(\frac{l\pi}{L_x} x\right) \cos\left(\frac{m\pi}{L_y} y\right) e^{j(\omega t - k_z z)} \end{array} \right)$$

For this problem, only include $l=0$, and the sum

$$\text{reduces to } \left(\begin{array}{l} e^{j\omega t} \\ \sum_{m=0}^{\infty} \hat{A}_{0m} \cos\left(\frac{m\pi}{L_y} y\right) e^{j(\omega t - k_z z)} \end{array} \right)$$

$$k_z = \pm \sqrt{k^2 - \left(\frac{m\pi}{L_y}\right)^2}, \quad \hat{k} = \frac{\omega}{c} - j\alpha$$

UNIVERSITY OF IDAHO Apply the BC's to determine the

\hat{A}_{0m} :

$$\hat{u}_z(y, z) = \frac{1}{\rho_0 \mu_0} \sum_{m=0}^{\infty} (k_z k_e) \hat{A}_{0m} \cos\left(\frac{m\pi}{L} y\right) e^{-jk_z z}$$

At $z=0$:

$$u_0 f(y) = \sum_{m=0}^{\infty} \frac{k_z}{\rho_0 \omega} \hat{A}_{0m} \cos\left(\frac{m\pi}{L} y\right) \neq \text{!}$$

Use the orthogonality of $\cos\left(\frac{m\pi}{L} y\right)$

2/1/4



$$\int_0^{ly} \cos\left(\frac{m\pi}{ly}y\right) \cos\left(\frac{q\pi}{ly}y\right) dy = \begin{cases} \frac{ly}{2}; & m=q \\ 0; & m \neq q \end{cases}$$

Then

~~$$\hat{A}_{0m} = \frac{\rho_0 \omega l_0}{k_2} \frac{1}{ly} \int_0^{ly} \cos\left(\frac{m\pi}{ly}y\right) f(y) dy$$~~

$$\hat{A}_{0m} = \begin{cases} \frac{\rho_0 \omega l_0}{k_2} \frac{1}{ly} \int_0^{ly} f(y) dy, & m=0 \\ \frac{\rho_0 \omega l_0}{k_2} \frac{2}{ly} \int_0^{ly} \cos\left(\frac{m\pi}{ly}y\right) f(y) dy, & m>0 \end{cases}$$

University of Idaho So now, we know the exact

acoustic pressure $j(\omega t - k_2 z)$

$$\hat{p}(x, z) e^{j\omega t} = \sum_{n=0}^{\infty} \hat{A}_{0n} \cos\left(\frac{n\pi}{L_y} y\right) e^{j(\omega t - k_2 z)}$$

$$k_z = \pm \sqrt{\hat{k}^2 - \left(\frac{n\pi}{L_y}\right)^2}$$

Axial wavenumber, let $\hat{k} = \frac{\omega}{c}$, $\alpha = 0$

$$k_z = \begin{cases} \pm \sqrt{\left(\frac{\omega}{c}\right)^2 - \left(\frac{n\pi}{L_y}\right)^2} & ; \frac{\omega}{c} > \frac{n\pi}{L_y} ; \text{Propagating} \end{cases}$$

$$k_z = \begin{cases} \pm j \sqrt{\left(\frac{n\pi}{L_y}\right)^2 - \left(\frac{\omega}{c}\right)^2} & ; \frac{\omega}{c} < \frac{n\pi}{L_y} ; \text{Evanescent} \end{cases}$$

— choose this root for decaying amplitude.

UNID University of Idaho for a fixed frequency ω ; define

$$M = \text{INT} \left(\frac{L_1 \omega}{\pi c} \right) \quad \text{MaxCAP INT} = \text{floor}$$

Then, The acoustic pressure can be split into two parts

$$p(y, z) e^{j\omega t} = \sum_{m=0}^M \hat{A}_{0m} \cos \left(\frac{m\pi}{L_1} y \right) e^{j(\omega t - k_2 z)} \quad \uparrow \quad k_2 = \sqrt{\left(\frac{m\pi}{L_1} \right)^2 - \left(\frac{\omega}{c} \right)^2}$$

propagating \rightarrow $\sum_{m=0}^{\infty} \hat{A}_{0m} \cos \left(\frac{m\pi}{L_1} y \right) e^{-\delta z}$ just

evanescent \rightarrow $\sum_{m=M+1}^{\infty} \hat{A}_{0m} \cos \left(\frac{m\pi}{L_1} y \right) e^{-\delta z}$ just $\delta = \sqrt{\left(\frac{m\pi}{L_1} \right)^2 - \left(\frac{\omega}{c} \right)^2}$

UNID University of Idaho The highest frequency at which
 Only the plane wave will propagate can be found from

$$\frac{\omega_c}{c} = \frac{1 \cdot \pi}{L_y}$$

$\omega_c =$ Cutoff Frequency

$$\omega_c = \frac{\pi}{L_y} c \Rightarrow$$

$$2 \pi f_c = \frac{\pi}{L_y} c$$

$$f_c = \frac{c}{2L_y}$$

$$L_y = \frac{1}{2} \frac{c}{f_c} = \frac{1}{2} \lambda$$

↙

Upper limit on waveguide

size for which only plane

wave will propagate at a freq
 f_c .